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Lattice dynamics and the NaCl-CsCl transformation in RbI

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It has been suggested (1-3) that the spontaneous transformation of RbI from the NaCl to the CsCl structure at high pressure occurs via a structural instability caused by vanishing of the frequency of the transverse acoustic lattice vibrational mode at the center of the $\langle 100 \rangle$ zone face, at a pressure near that of the transformation. This problem is discussed and very much extended by HARDY and KARO.⁽³⁾ Their results indicate that, for appropriate choice of interionic interactions, as the solid is compressed toward the transition conditions the acoustic branches of the lattice vibrational spectrum are strongly shifted to lower energies while the optical branches are shifted slightly to higher energies as shown in Fig. 2 of HARDY and KARO.⁽³⁾ The purpose of this note is to point out that the results quoted are strongly inconsistent with the observed values (4-6) of thermal expansion of RbI if the energy shifts calculated are linear over the pressure range to the transition.

Taking the familiar Quasi harmonic oscillator⁽⁷⁾ point of view the lattice excitations are considered to be those of a set of nearly harmonic oscillators which exhibit a volume dependence of their characteristic energies. This approximation leads to the result for "Grüneisen's Gamma":

$$\gamma_G \equiv \frac{\alpha B_T V}{C_v} = \frac{\sum_i \gamma_i C_{vi}}{\sum_i C_{vi}}$$

in which α is the thermal expansion coefficient, B_T is the Bulk Modulus C_v/V is the heat capacity per unit volume,

$$\gamma_i \equiv \frac{d \ln \nu_i}{d \ln V}$$

a so called "mode gamma", ν_i is any lattice vibrational frequency and C_{vi} is the Einstein heat capacity of the *i*'th mode at the temperature of observation. The summation is over all oscillators making up the crystal vibrational spectrum. At high temperature each oscillator is classically excited and its heat capacity equals the Boltzmann constant k, whence

$$\gamma_{\infty} = \frac{\sum_{i} \gamma_{i}}{3N}$$

the simple average of the mode gammas. At intermediate temperatures one sees that the mode gammas are weighted by their associated Einstein heat capacities at the observation temperature. Let us perform a qualitative analysis, estimating mode gammas from Fig. 2 of HARDY and KARO.

Compression data from Slater yields about 4 percent change in volume to the transition pressure. The average shift of the collected acoustic mode peaks is about 15 percent yielding an average acoustic mode gamma -4. Similarly, the average optical mode gamma is about +0.7 yielding a high temperature limiting value of $\Gamma_G = -3.3$ decreasing algebraically with decreasing temperatures. Since in equation (1) B_T , V and C_v must be positive, the implication of the negative gamma is that the thermal expansion coefficient have negative values at all temperatures. The thermal expansion coefficient of RbI has been measured over a wide temperature range.⁽⁴⁻⁶⁾ Its value becomes negative at temperatures below 10°K but even at 5°K the value of γ estimated from the expansion data is about -0.2×10^{-4} , still less negative than value -4 suggested by the lattice dynamical calculations. (The low temperature limit of γ calculated from measurements of the pressure dependence of the elastic constants of RbI is positive.)

BARRON's suggestion⁽⁷⁾ that the thermal expansion of solids is a quantity which yields information about the microscopic interactions in the solid is relevant here. In the present case, it seems that the requirement to account for the experimental values of the thermal expansion *and* to find a mode frequency which will vanish to permit spontaneous NaCl–CsCl transition will provide a challenge for the lattice dynamicists. We feel at present that the mode instability hypothesis probably is incorrect as the explanation of the spontaneous nature of the transformation on RbI.

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